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Secondary Guest-Driven Nucleation of N₂O Clathrate Hydrate in Amorphous Ice Mixture under Ultrahigh Vacuum

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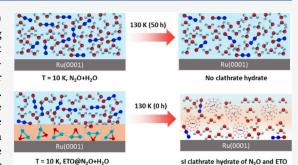
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ABSTRACT: We report the first observation of nitrous oxide (N_2O) clathrate hydrate (CH) under ultrahigh vacuum (UHV) conditions using infrared spectroscopy. While nucleation of CHs from water and guest molecules either spontaneously or in the presence of promoters is wellestablished under high-pressure conditions, their occurrence under cryogenic UHV remains less understood, particularly in multiguest systems. In this study, binary and ternary ice mixtures of N_2O , ethylene oxide (ETO), dimethyl ether (DME), tetrahydrofuran (THF), and H_2O were investigated via thermal annealing at 130 K. The binary N_2O-H_2O system showed no signs of hydrate formation after 50 h of annealing, whereas the addition of ETO, DME, or THF led to immediate CH formation under identical conditions. This demonstrates a clear promoter-assisted



heterogeneous nucleation mechanism. In situ reflection absorption infrared (RAIR) spectroscopy confirms that ETO facilitates the formation of structure I hydrates, incorporating both ETO and N_2O in the cages, while DME and THF promote larger-cage CHs. Density functional theory calculations support experimental RAIR results, which predict stronger host—guest interactions and blue shifts for small cages and weaker shifts for large cages. These findings offer new mechanistic insights into multiguest hydrate nucleation and establish N_2O CH formation under UHV as a relevant process in cryogenic and astrochemical environments.

Titrous oxide (N_2O) is a potent greenhouse gas with a global warming potential approximately 300 times greater than that of carbon dioxide (CO₂) over a 100-year time scale. Beyond its climatic impact, N2O contributes to stratospheric ozone depletion through photodissociation. According to the U.S. Environmental Protection Agency (EPA), anthropogenic sources accounted for about 5% of total greenhouse gas emissions in the United States in 2013. Interestingly, N2O is not confined to Earth; it has also been detected in interstellar regions such as Orion KL and Sagittarius B2, indicating that nitrogen-oxygen bond formation occurs under extreme conditions of the interstellar medium (ISM).^{4,5} Recent observations from the James Webb Space Telescope (JWST) suggest the presence of condensedphase N₂O in protostellar environments, further underscoring its astrochemical relevance.

Considering the presence of N_2O in terrestrial and astrophysical environments, understanding its potential for physical trapping and chemical transformation in icy matrices is of significant interest. One mechanism by which volatile molecules such as N_2O may be sequestered is through clathrate hydrates (CHs). CHs are crystalline solids made of hydrogen-bonded water molecules that form polyhedral cages that encapsulate guest species such as CH_4 , CO_2 , ethane, etc.^{7,8} These structures, which typically adopt structure I (sI), structure II (sII), or structure H (sH), naturally form under

high-pressure and low-temperature conditions found in deep ocean sediments and planetary ices. ^{9,10} CHs are believed to play a major role in cometary outgassing events. The observed outgassing patterns in the Northern Hemisphere suggest that comet 67P's nucleus contains a combination of crystalline ice, CHs, and other ices, contradicting the hypothesis that its volatiles were primarily trapped in an amorphous ice phase. ¹¹ Recent laboratory experiments have demonstrated that CHs can also form under UHV and cryogenic conditions, broadening the potential relevance of CHs to astrophysical ices. ^{12–14} For instance, thin-film clathrates of THF have been successfully synthesized and structurally confirmed via electron diffraction. ¹⁵

Despite growing evidence for clathrate formation under nontraditional conditions, the fundamental process of nucleation remains poorly understood. Classical nucleation theory (CNT) suggests that hydrate formation begins once a critical cluster of host and guest molecules is formed.¹⁶ In

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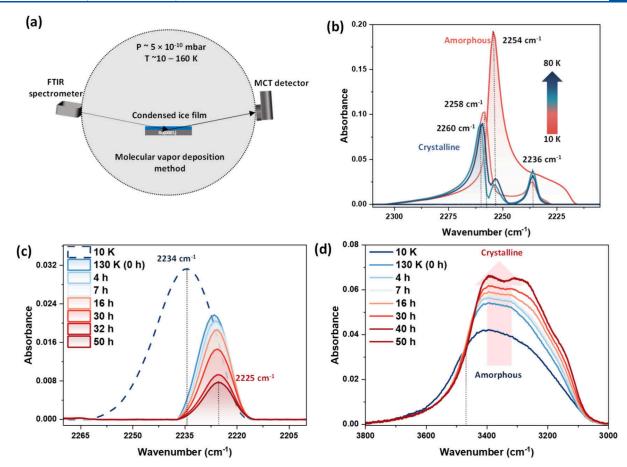


Figure 1. (a) A simplified schematic of the UHV instrument used for this work. (b) Temperature-dependent RAIR spectra of 150 ML of N_2O ice in the N-N-O asymmetric stretching region. Temperature- and time-dependent RAIR spectra of 300 ML of $N_2O + H_2O$ ice in the N-N-O asymmetric stretching region (c) and in the O-H stretching region (d).

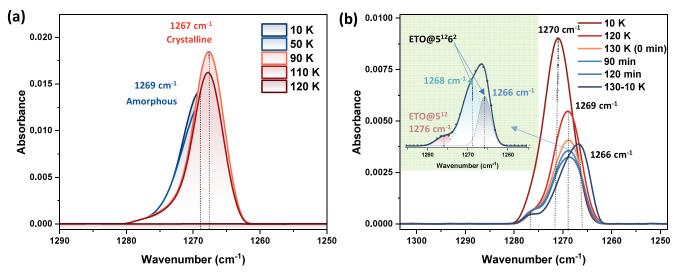


Figure 2. (a) Temperature-dependent RAIR spectra of ETO ice in the ring-breathing region. (b) RAIR spectra of ETO + H_2O ice in the ring-breathing region. The spectrum, acquired after annealing at 130 K for 2 h and subsequently cooling to 10 K, is deconvoluted into three distinct peaks located at 1276 cm⁻¹ (red), 1268 cm⁻¹ (sky blue), and 1266 cm⁻¹ (blue).

contrast, simulations propose a two-step mechanism involving the formation of an amorphous precursor that later crystallizes into a stable hydrate structure. Our prior work using amorphous mixtures of water and various guest molecules under UHV revealed that most CHs formed under these conditions are metastable, often reverting to cubic or

hexagonal ice. $^{20-22}$ Among the systems tested, only tetrahydrofuran (THF, C_4H_8O) consistently formed stable CHs, and these findings were attributed primarily to homogeneous nucleation processes. 15,23

Building upon these observations, we turned our attention to N_2O , a known sI clathrate former under high-pressure

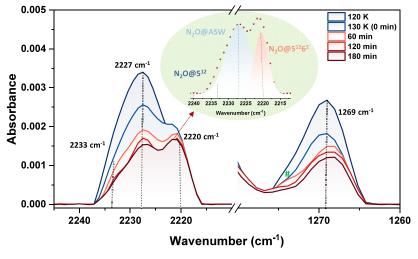


Figure 3. RAIR spectra of ETO@ $N_2O + H_2O$ showing the N-N-O asymmetric stretching region of N_2O and the ring-breathing mode of ETO. After 180 min, the N-N-O stretching band is deconvoluted into three components, with peaks centered at 2233 cm⁻¹ (deep blue), 2227 cm⁻¹ (light blue), and 2220 cm⁻¹ (red). The peak labeled as # in the spectra in the ring breathing mode is attributed to the ETO occupied in the small cage.

Table 1. Calculated Infrared Frequencies of N_2O and ETO in the N-N-O Asymmetric Stretching Region and the Ring-Breathing Region

| System | N_2O | $N_2O@5^{12}$ | $N_2O@5^{12}6^2$ | N ₂ O@5 ¹² 6 ⁴ |
|-------------------------------|---------|---------------------|------------------------------------|---|
| Frequency (cm ⁻¹) | 2267.93 | 2286.69 | 2272.89 | 2269.74 |
| | | | | |
| System | ETO | ETO@5 ¹² | ETO@5 ¹² 6 ² | ETO@5 ¹² 6 ⁴ |

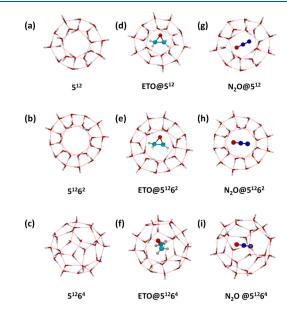


Figure 4. Optimized geometries of (a-c) empty 5^{12} , $5^{12}6^2$, and $5^{12}6^4$ CH cages; (d-f) ETO inside the 5^{12} , $5^{12}6^2$, and $5^{12}6^4$ CH cages; (g-i) nitrous oxide encaged in the 5^{12} , $5^{12}6^2$, and $5^{12}6^4$ CH cages.

conditions, as supported by X-ray diffraction and Raman data. We aimed to investigate whether N_2O could form CHs from amorphous N_2O-H_2O ice mixtures under UHV and cryogenic conditions, which more closely resemble interstellar ice environments. Surprisingly, N_2O alone failed to nucleate a clathrate phase in contrast to other guest molecules. To address this, we introduced ethylene oxide

(ETO/ C_2H_4O), a molecule previously shown to nucleate sI CHs under vacuum conditions, as a coguest. PRemarkably, as ETO initiated CH formation, N_2O got incorporated into the clathrate lattice, suggesting a conucleation mechanism or surface-mediated heterogeneous nucleation. This mechanism is further supported by additional experiments involving other known hydrate formers such as THF and dimethyl ether (DME/C_2H_6O) .

This work presents the first IR spectroscopic evidence of the existence of N_2O CH under UHV conditions. The ability of N_2O to form mixed CHs with ETO extends the known range of CH-forming systems and offers important implications for astrochemistry. In particular, such mixed ices could serve as precursors for complex organic molecules, including amino acids and nucleobases, when exposed to ionizing radiation in space.³⁰ This study provides new insights into the physical chemistry of clathrate nucleation and potential pathways for molecular evolution in icy astrophysical environments.

All experiments were conducted using a custom-built ultrahigh vacuum (UHV) chamber, as described in detail in our previous study (Figure 1a). Thin ice films of N_2O , H_2O , ETO, DME, and THF were prepared by background vapor deposition. The ice thickness was estimated based on the ion gauge sensitivity factor, deposition time, and background pressure. Quantum chemical calculations for free and encapsulated N_2O and ETO were performed using the B3LYP³² functional with the 6-311++G(d,p)³³ basis set. Additional methodological details are provided in the Supporting Information.

At the outset, 150 ML of N_2O ice were deposited on a Ru(0001) substrate at 10 K. Figure 1b and Figure S1 present the temperature-dependent RAIR spectra of the deposited N_2O ice in the N-N-O asymmetric and symmetric stretching region. At 10 K, a sharp peak at 2254 cm⁻¹ and a broad feature in the lower wavenumber region (2245–2215 cm⁻¹) were observed, indicating that the N_2O ice is in an amorphous state. Since N_2O is structurally similar to CO_2 , it exhibits two distinct peaks for the ν_3 stretching mode, corresponding to the transverse optical (TO) and longitudinal optical (LO) modes. The higher wavenumber feature (2254 cm⁻¹) corresponds to the LO mode, while the lower wavenumber broad peak

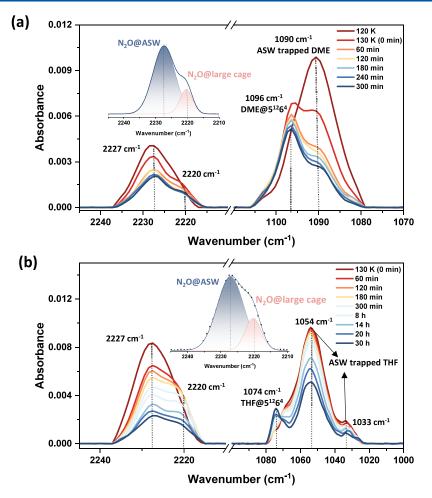


Figure 5. (a) RAIR spectra of DME@ $N_2O + H_2O$ in the N-N-O asymmetric stretching region of N_2O and the C-O-C asymmetric stretching region of DME. After 300 min, the N-N-O stretching band is deconvoluted into two components, centered at 2227 cm⁻¹ (highlighted in deep blue) and 2220 cm⁻¹ (highlighted in red). (b) RAIR spectra of THF@ $N_2O + H_2O$ in the N-N-O asymmetric stretching region of N_2O and the C-O asymmetric stretching region of THF. Similarly, the spectrum after 300 min shows the N-N-O stretching band resolved into two components at 2227 cm⁻¹ (deep blue) and 2220 cm⁻¹ (red).

represents the TO mode. $^{34-37}$ The RAIRS measurement was performed at a grazing incidence angle of 82.75° , 31 which enhances the contribution from the LO mode, thus explaining the higher intensity of the 2254 cm $^{-1}$ peak. We note that the experiments were performed on a Ru metal surface, where the surface dipole selection rule suppresses the TO mode, thereby making the LO mode more prominent. 38 Consistent with this substrate dependence, Cassidy et al. 39 reported two distinct infrared spectra for the N–N–O asymmetric stretch of N₂O on silica versus Cu: the silica support retains TO activity, whereas the metallic Cu surface, subject to the same selection rule as Ru, leads to an LO-dominated profile. 39

Upon annealing the N_2O ice to its desorption temperature (~80 K), the amorphous ice undergoes crystallization, likely forming cubic crystalline ice in two stages. During this process, the 2254 cm⁻¹ peak splits into two components (2258 and 2254 cm⁻¹), and its intensity decreases. The broad peak also sharpens, evolving into a distinct peak at 2236 cm⁻¹. This decrease in intensity is attributed to dipole alignment during crystallization, which reduces the dipole moment fluctuations and hence the infrared absorption. Similar behavior has been previously reported using various spectroscopic techniques. Property of the symmetric stretching region (Figure S1), a broad peak at 1296 cm⁻¹ indicates the amorphous phase and a

sharp peak at $1299~{\rm cm}^{-1}$ indicates the crystalline phase. We note here that $150~{\rm ML}$ of N_2O ice RAIR spectra are also given in our earlier paper, ⁴⁴ Supporting Information, but a detailed analysis of spectral features has not been done.

To investigate the formation of N2O CH, 300 ML of a N2O:H2O ice mixture in a 1:5 ratio was codeposited and thermally annealed to 130 K, where it was held for 50 h. During the initial warming from 10 to 130 K, a red shift was observed, accompanied by a decrease in the peak intensity of the N-N-O asymmetric stretch (Figure 1c, Figure S2), which is attributed to the desorption of N2O from the ice matrix. Upon annealing at 130 K, the band intensity decreases with no further frequency change, indicating no new interaction, only a declining N₂O/H₂O ratio. Thus, once the frequency stabilizes, the number of H₂O molecules coordinated to N₂O appears to be constant. Additionally, in the symmetric stretching region, the same trend was observed, where we only saw a decrease in the peak position at 1299 cm⁻¹ (Figure S3). We note here that the available data do not allow us to distinguish specific N₂O-H₂O conformations, so we refrain from making a conformer assignment. Additional temperature-dependent experiments with 1:10 and 1:20 of N2O:H2O (Figure S4) ratios showed similar behavior, peak red-shifting and decreasing intensity, likely due to enhanced hydrogen bonding between the N₂O

and surrounding water molecules.³⁵ Unlike our earlier work, N₂O does not exhibit facile homogeneous nucleation of CHs. In the O-H stretching region (Figure 1d), sharpening of the peaks was observed, consistent with the formation of crystalline water ice. 45 At low N₂O concentrations in H₂O ice, the O-H stretching region is only weakly perturbed. Increasing the mixing ratio to 1:1 (N₂O:H₂O) produces a prominent dangling O-H feature at 3652 cm⁻¹ (Figure S5). A weaker hump at 3491 cm⁻¹ is also observed, attributable to N₂O-H₂O interactions, although its specific origin remains unclear. Upon warming to 100 K, both peaks vanish, indicating compaction of the ice matrix and reduced N2O content. These experiments show that N2O CH does not form spontaneously from an N₂O-H₂O ice mixture under UHV conditions. Consequently, we focused on the formation of sI ETO CH, which was later used as a promoter for N₂O CH formation.

To begin, pure ETO ice was investigated under UHV conditions. Figure 2a shows the temperature-dependent RAIR spectra of ETO in the ring-breathing region (v_3). At 10 K, vapor-deposited ETO formed an amorphous ice, evident from a broad peak centered at 1269 cm $^{-1}$. Upon annealing to 90 K, crystallization occurred, as indicated by the emergence of a sharp peak at 1267 cm $^{-1}$. In the CH $_2$ bending region, the initially broad band sharpens and resolves into two peaks at 1466 and 1460 cm $^{-1}$ (Figure S6a). Upon further warming to 120 K, ETO completely desorbs from the surface.

To study the formation of ETO CH, approximately 300 ML of ETO– $\rm H_2O$ (1:5) vapor mixture was codeposited and annealed to 130 K (Figure 2b). It is well-established that ETO forms sI CH under both high-pressure and vacuum conditions, and its characteristic RAIR spectra are also well-documented. In particular, the ring-breathing mode of ETO in sI CHs typically exhibits a doublet with a 2 cm⁻¹ splitting, attributed to ETO molecules occupying large cages in two distinct orientations. A smaller hump, ~ 13 cm⁻¹ blue-shifted from the main peak, has been assigned to ETO trapped in smaller cages. Although this splitting is not always fully resolved at elevated temperatures, it becomes clearer at lower temperatures due to reduced thermal broadening.

In our experiment, after annealing to 130 K, we observed a broadened peak with a small shoulder at around 1276 cm⁻¹ (Figure 2b), suggesting the formation of clathrate structures. Upon cooling the sample back down to 10 K, a clearer spectral splitting emerged with increased peak intensity, indicative of more ordered structural arrangements. Deconvolution of the spectrum revealed two sharp peaks at 1268 cm⁻¹ and 1266 cm⁻¹, along with a smaller intensity feature at 1276 cm⁻¹ (Figure 2b). This splitting is possible due to the different orientation of ETO inside sI large cages.⁴⁷ The peak positions are consistent with those in earlier reports. Bertie and Othen⁴⁸ observed features at 1268 and 1270 cm⁻¹, while Fleyfel and Devlin⁴⁷ reported 1266 and 1268 cm⁻¹ for ETO in a 5¹²6² cage. Taken together, these observations suggest the facile formation of ETO CH from the amorphous ice mixture at 130 K. In the CH₂ bending region (Figure S6b), no definitive changes were observed due to very low intensity. Concurrently, crystallization of amorphous solid water (ASW) was evident, as indicated by the sharpening of features in the O-H stretching region (Figure S7). While ETO CH formation under vacuum has been previously reported, 48,53 its formation from in ASW under UHV conditions has not been demonstrated. One motivation for studying ETO CH in this manner is to examine

whether the surface of ETO CH can facilitate the nucleation of N₂O CH.

In the first attempt, we failed to create N2O CH by thermal annealing the N₂O-H₂O ice mixture. Later, we took a N₂O-H₂O ice mixture on top of ETO ice and, through thermal annealing, created a mixture of ETO-N2O-H2O ice at 130 K (Figure 3). When ETO started to form sI CH, a clear splitting was observed in the N-N-O asymmetric stretching region. The peak at 2227 cm⁻¹ was split to make a new peak at 2220 cm⁻¹, and a small hump was observed at 2233 cm⁻¹. The peak intensity at 2233 cm⁻¹ decreases relatively slowly (and remains nearly constant over 120-180 min), compared to the peak intensity of 2227 cm⁻¹ (Figure 3). This behavior indicates two N₂O populations, with the persistent component stabilized in well-defined sites consistent with N2O trapped in CH cages, whereas the faster-decreasing band reflects more weakly bound ASW-trapped N₂O that desorbs. The area under the two peaks (at 2233 cm⁻¹ and 2220 cm⁻¹) gave a ratio of 1:4.5, which is close to the population distribution in the small and large cages of sI CH. In the symmetric stretching region (Figure S8) two new peaks were observed at 1282 and 1292 cm⁻¹. It is worth noting that peak deconvolution can not be done for symmetric stretch due to a poor signal-to-noise ratio. Additionally, in the ring-breathing mode of ETO, a broadened peak at 1269 cm⁻¹ and a small hump at 1276 cm⁻¹ (Figure 3) were observed, consistent with previous observations of pure ETO sI CH.⁴⁸ A Raman spectroscopic study reported by Yang et al.⁵⁴ shows the difference between the large cage and small cage peak is 13 cm⁻¹ (for asymmetric stretch) and 8 cm⁻¹ (for symmetric stretch), respectively, which closely match with the current IR study. This proves the formation of sI N₂O CH in the presence of ETO CH.

The nucleation behavior of CH from N2O is notably complex. In the absence of ETO, prolonged thermal annealing over several hours failed to induce CH formation; instead, the ASW began to crystallize. This observation clearly indicates that the homogeneous nucleation of CH from an amorphous N₂O:H₂O ice mixture is not feasible without the presence of a third component. However, once the sI CH formation of ETO commenced in the ETO-N2O-H2O ice mixture, CH formation from N₂O was also observed. Two possible nucleation pathways can be considered: (i) simultaneous nucleation of ETO and N2O hydrates due to their homogeneous distribution in the ice matrix, or (ii) heterogeneous nucleation of N2O hydrate on the surface of the preformed ETO hydrate. Given the limited experimental technique, it remains challenging to distinguish between these mechanisms or to rule out the possibility that both are occurring concurrently.

To support the IR results, we carried outgas-phase DFT calculations for N₂O and ETO in sI and sII CH cages. The optimized geometries of the empty 5¹², 5¹²6², and 5¹²6⁴ CH cages, as well as those containing ETO and N₂O, are shown in Figure 4. The vibrational frequencies of the isolated ethylene oxide and nitrous oxide molecules, along with their frequencies when confined within the three different cages, are shown in Table 1. Free N₂O exhibits an asymmetric stretch at 2267.93 cm⁻¹, consistent with a prior report,⁵⁵ and ETO shows ring breathing mode at 1257.19 cm⁻¹. When these molecules are confined within the smaller 5¹² CH cage, they exhibit blue shifts of 18.76 cm⁻¹ and 15 cm⁻¹ for N₂O and ETO, respectively. However, when placed inside the comparatively larger 5¹²6² and 5¹²6⁴ CH cages, the N₂O molecule shows

lower blue shifts of 4.96 cm⁻¹ and 1.81 cm⁻¹, while the ETO molecule exhibits blue shifts of 5.51 cm⁻¹ and 1.98 cm⁻¹, respectively. The molecules exhibit a stronger blue shift when confined within smaller cages due to steric restrictions that distort molecular bonds and elevate vibrational frequencies. In contrast, larger cages provide greater spatial freedom and structural flexibility, minimizing steric effects and resulting in only slight changes in vibrational modes.⁵⁶ It is important to note that no chemical bond formation was observed between the water molecules of the CH cages and guest molecules (N2O and ETO). Though we observed distortion of the CH cages, we did not observe cage breakage while guest molecules were confined in them. Hydrogen bonding between the oxygen and hydrogen molecules forming the CH cages and weak interaction between the cage and guest molecules provide the stability to molecule-confined CH cages.⁵⁷

To further prove the heterogeneous nucleation of N₂O CH in amorphous ice, we have used two stable sII CH former molecules, DME and THF (Figure 5).⁵⁸ Both of the molecules are known to form sII large cage CH in UHV under cryogenic conditions. We deposited sequentially THF/DME and H_2O + N₂O ice mixture at 10 K and annealed the sample at 130 K for a few hours. Figure 5 shows IR spectra of DME@H₂O+N₂O (Figure 5a) and THF@H₂O+N₂O (Figure 5b) ice mixture in the N-N-O and C-O asymmetric stretch. Previous studies showed that THF and DME form sII CH where the 5¹²6⁴ cage is occupied by the THF and DME with a distinct assigned IR peak 1074 cm⁻¹ and 1096 cm⁻¹. ^{22,53,59} The peaks at 1090 cm⁻¹ and 1033/1054 cm⁻¹ are attributed to DME and THF molecules trapped within ASW. In this system, during the formation of THF/DME CHs in the N-N-O asymmetric stretching region a new peak arises at 2220 cm⁻¹, which can be assigned to the large cages. Interestingly, we have not observed any peak around 2233 cm⁻¹, which indicates the absence of small cages. Here, it is difficult to assign whether this assignment occurs due to the large cages of sI or sII. It is worth noting that N2O is not going to preexisting vacant small cages created by THF/DME. These additional experiments demonstrate that, without any incubation time at 130 K, N2O can spontaneously form a CH only in the presence of supporting guest molecules such as ETO, DME, or THF, which facilitate nucleation.

In conclusion, we have demonstrated the formation of N₂O and ETO CHs under UHV and cryogenic conditions. While the formation of ETO CHs in high vacuum was previously reported by the Devlin group, 28,58 to the best of our knowledge, this is the first report of ETO CH formation under such conditions. Notably, we also present the first IR spectra of N₂O CHs. A key conclusion from this study is that homogeneous nucleation of N2O CH does not occur readily from an amorphous N2O-H2O ice mixture, even upon annealing to 130 K. However, when codeposited with a secondary guest molecule known to spontaneously form clathrates, such as ETO, DME, or THF, N2O readily forms CHs upon annealing. In the presence of ETO, N₂O forms sI CHs, while in the presence of DME or THF, N2O appears to occupy larger cages; however, due to limited spectroscopic evidence, we could not conclusively assign these to either sI or sII CHs. To support our experimental observations, DFT calculations were performed for free and encaged N2O and ETO molecules within 5^{12} , $5^{12}6^2$, and $5^{12}6^4$ cages. The calculated IR spectra indicate that molecules confined in

smaller cages experience greater steric hindrance, resulting in larger blue shifts compared to those in larger cages.

ASSOCIATED CONTENT

Solution Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acs.jpclett.5c02778.

Experimental section (including experimental setup, materials and reagents, and sample preparation) and temperature and time-dependent RAIR spectra (PDF) Transparent Peer Review report available (PDF)

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Author Contributions

*B.K.M. and S.C. contributed equally. B.K.M. designed the experiments. B.K.M., S.C., S.S.B., and T.J.M. performed the experiments. B.K.M. and S.C. analyzed the results. T.P. and R.K. supervised the progress. S.K.N. and S.S.R.K.C.Y. performed the quantum chemical calculations. B.K.M. wrote the first draft of the manuscript. The final version of the manuscript was prepared with the contributions of all authors.

Notes

The authors declare no competing financial interest.

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